

Shock System Stability and Inlet-Isolator Flow Path Control with Scale-Resolved Simulations

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Objectives

- Identify causes of low-frequency streamwise shock-train oscillations in rectangular cross-section geometry ducts.
- Isolate upstream-propagating perturbations and assess whether their characteristic velocity supports low-frequency oscillation.
- Explore propagation pathways and modes of upstream travel.

Problem Definition

Experimental reference (*Benton et al. 2024 [1]*)

Behavior:

- Unsteady shock streamwise meandering known as shock-train unsteadiness

Repercussions

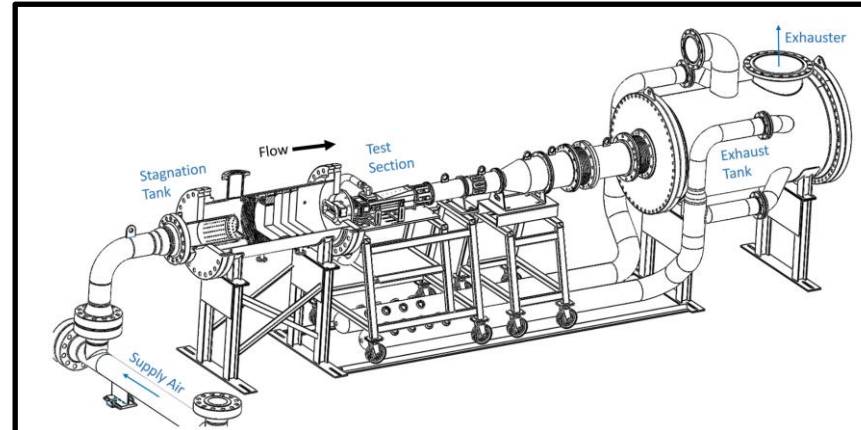
- Unsteady loads on vehicle
- Unsteady pressure recovery, combustion, and thrust
- Propensity for unstart

Possible causes:

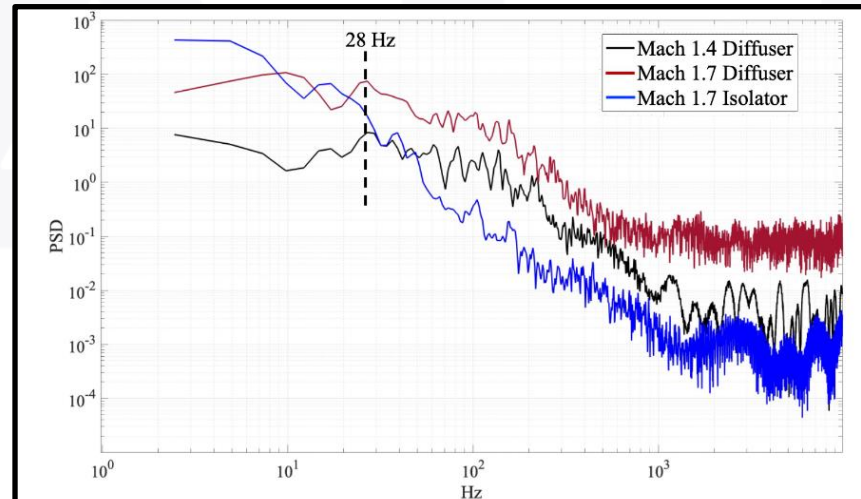
- Innate shock-boundary layer interaction (SBLI) instabilities [2]
- Upstream turbulent boundary layer super-structures [3]
- Downstream perturbations [4]

Experimental Conditions [5]

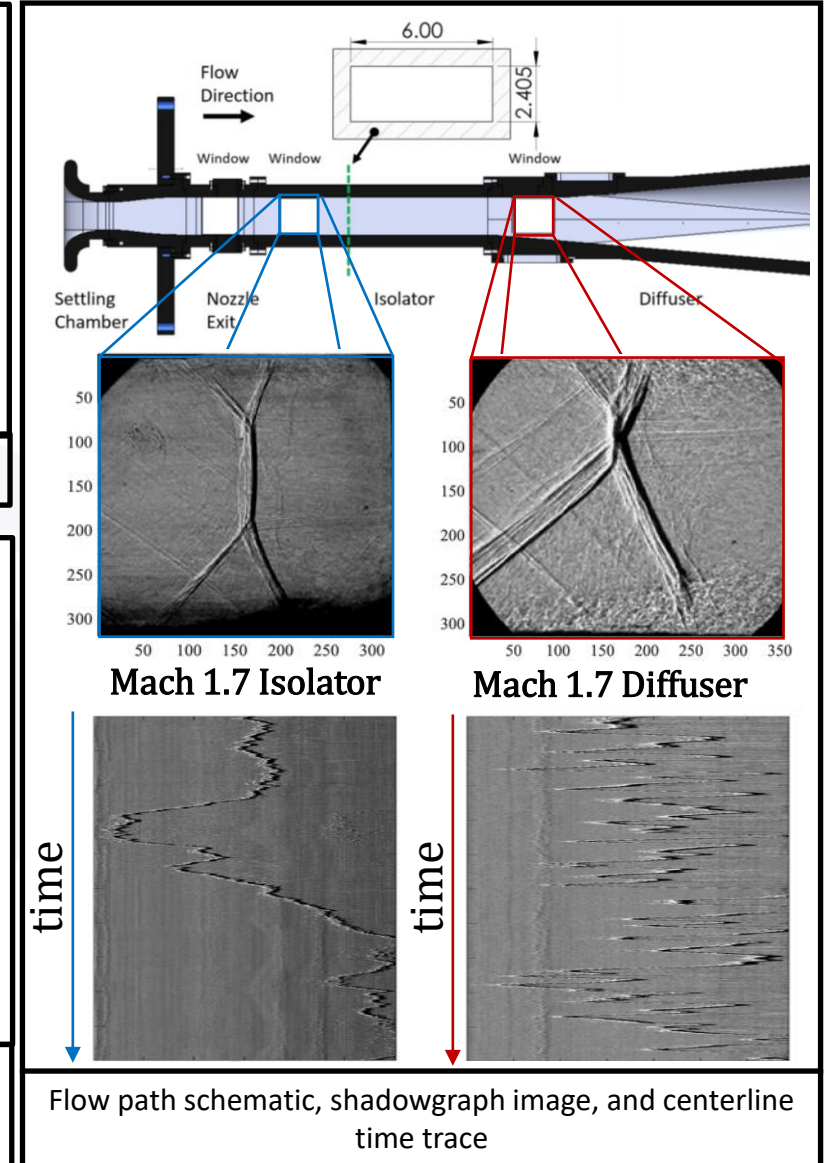
- $Mach = 1.4, 1.7$
- $P_O = 45$ psia
- $T_O = 491$ °R



Direct-Connect Transonic Internal Flow Facility

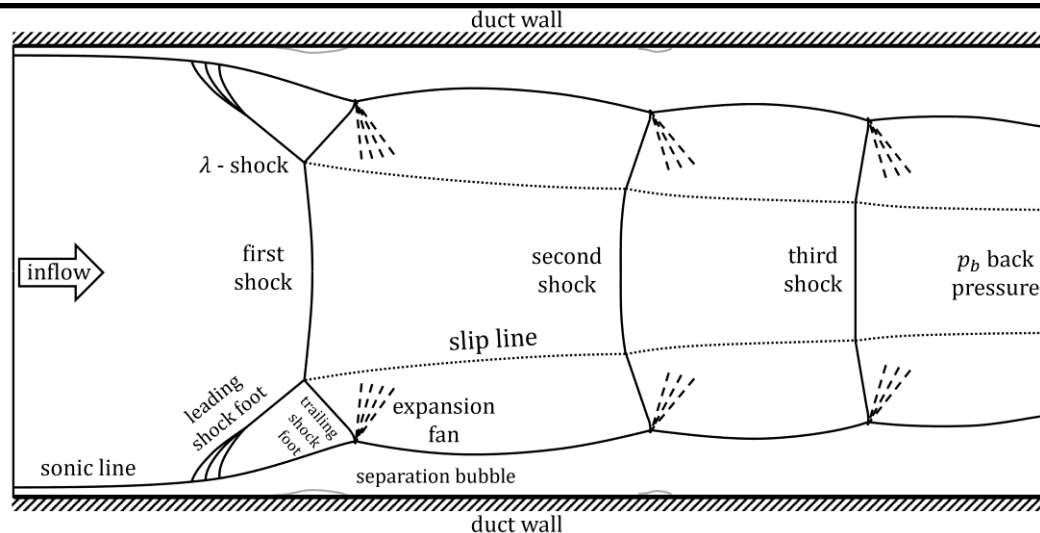


Low-frequency peaks from shadowgraph time traces of first normal shock.

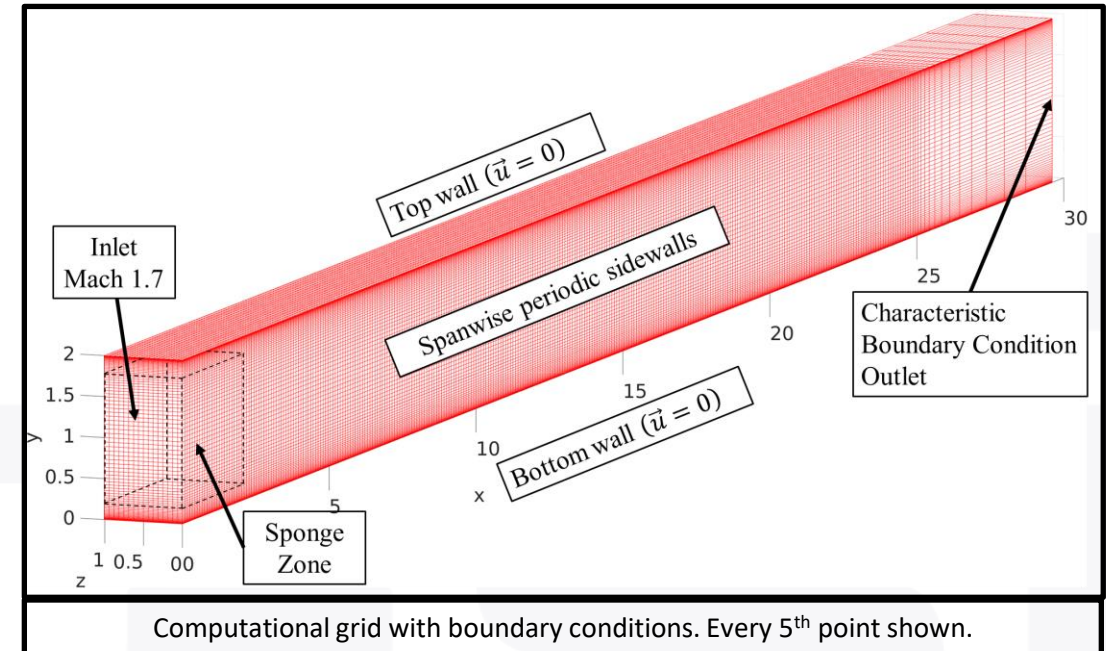


Computational Approach: Large Eddy Simulations (LES)

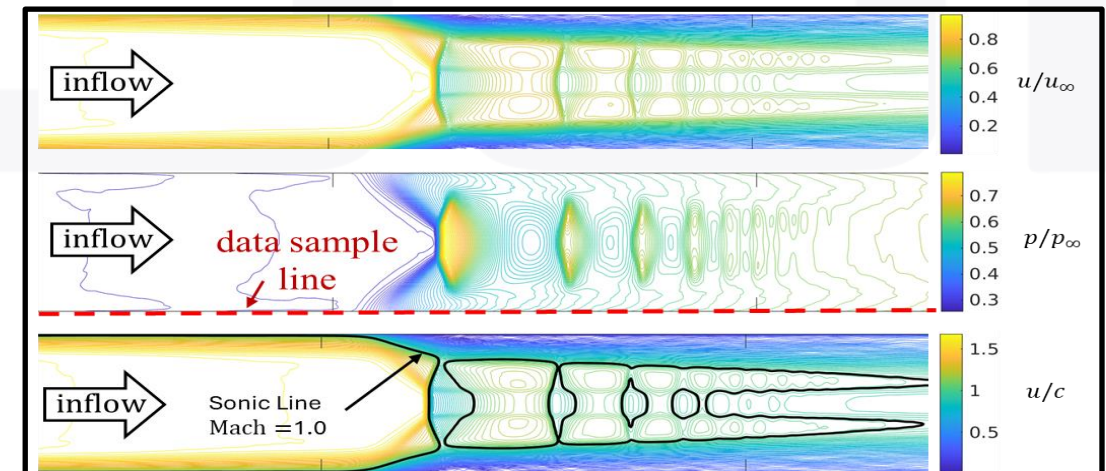
- High-fidelity computational fluid dynamics used to simulate the isolator shock-train.
- $Mach = 1.7$, $Re_h = 50,000$ (h = channel half height)
- Computational Domain
 - $(\hat{x}, \hat{y}, \hat{z}) = (29.6061, 2.0, 1.0)$
 - $(\hat{i}, \hat{j}, \hat{k}) = (1241, 321, 65)$
- Spatial scheme: 6th order compact, switching to 3rd order Roe near shock discontinuities.
- Spatial filtering (oscillation damping): 8th order implicit.
- Time integration: 3rd order Runge-Kutta.
- Digital filtering (turbulent inflow). No sub grid model



Schematic of shock-train structure with shock-boundary layer interactions (SBLI)



Computational grid with boundary conditions. Every 5th point shown.



Mean flow profiles from LES. Top, u-velocity. Middle, pressure. Bottom, Mach number

Doak's Momentum Potential Theory: Acoustics

Objective: Identify acoustic content

Tool: Doak decomposition separates acoustics from hydrodynamic and thermal content.

Process:

A Helmholtz Decomposition separates mass flux (ρu_i) into solenoidal (B_i) and irrotational ($-\nabla\psi$) components.

Continuity equation invoked.

The irrotational ($-\nabla\psi$) component is further decomposed into acoustic ($-\nabla\psi_A$) and thermal ($-\nabla\psi_T$) fields which are solved via means of three Poisson equations (in practice only (1) and (2) must be solved)

P.E. Doak 1989 [4]

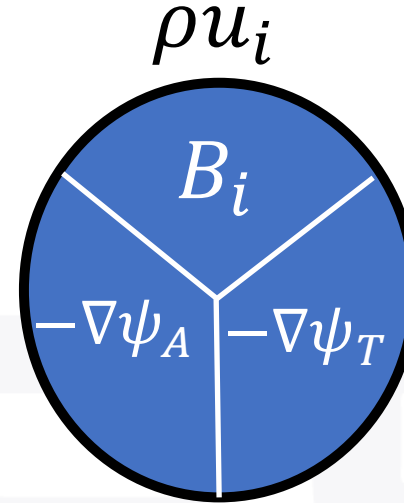
$$\rho u_i = B_i - \nabla\psi$$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho u_i) = 0$$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (B_i - \nabla\psi) = 0$$

$$\nabla \cdot B = \nabla \cdot \bar{B} + \nabla \cdot B'_t = 0$$

$$\nabla\psi = \nabla \cdot \bar{\psi}_t + \nabla\psi';$$



$$\nabla^2 \psi' = \frac{\partial \rho'}{\partial t} \quad (1)$$

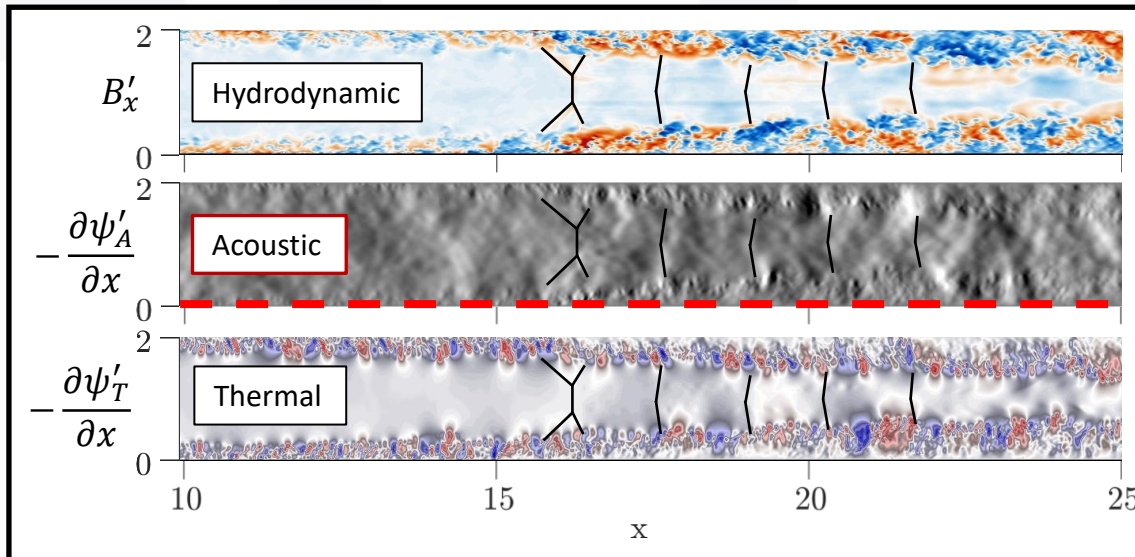
$$\psi' = \psi'_A + \psi'_T$$

$$\nabla^2 \psi'_A = \frac{1}{c^2} \frac{\partial p'}{\partial t} \quad (2)$$

$$\nabla^2 \psi'_T = \frac{\partial \rho}{\partial S} \frac{\partial S'}{\partial t} \quad (3)$$

Where S = entropy

Poisson equations



Space-Time Analysis

Purpose: To visualize direction and speed of acoustic signal motion near isolator wall

Procedure:

Sample streamwise line near bottom wall and plot over time. Fluctuations only.

Significance: $v = 1/\Delta t/\Delta x$

Positive-slope = downstream | Low-slope = fast

Negative-slope = upstream | High-slope = slow

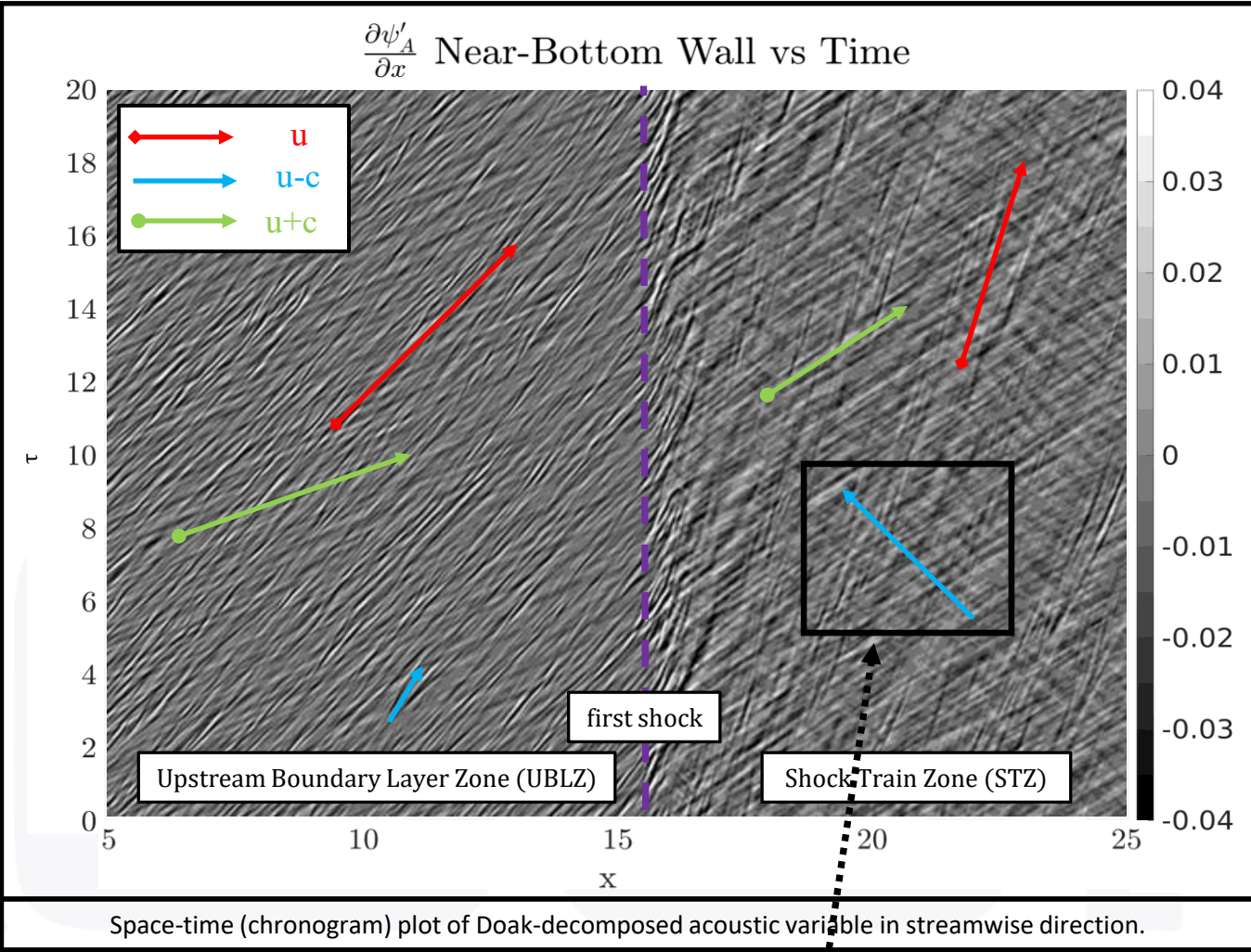
Three families of sloped lines per zone.

Upstream boundary layer zone has no upstream fluctuations

Shock-train zone has one upstream fluctuation

Can calculate velocity of individual structures.

However, only measuring one structure is undesirable, and measuring many lines is tedious.



| | | <i>u</i> | <i>u + c</i> | <i>u - c</i> | $c^+ = (u + c) - u$ | $c^- = u - (u - c)$ | c_{avg} |
|-------------|--|----------|--------------|--------------|---------------------|---------------------|-----------|
| UBLZ | $-\frac{\partial \psi'_A}{\partial x}$ | 0.720 | 1.836 | 0.430 | 1.116 | 0.290 | 0.703 |
| STZ | $-\frac{\partial \psi'_A}{\partial x}$ | 0.197 | 1.180 | -0.727 | 0.983 | 0.924 | 0.954 |

Upstream-running acoustics

Frequency-Wavenumber analysis

Purpose: To measure velocity of an ensemble of acoustic structures

Procedure:

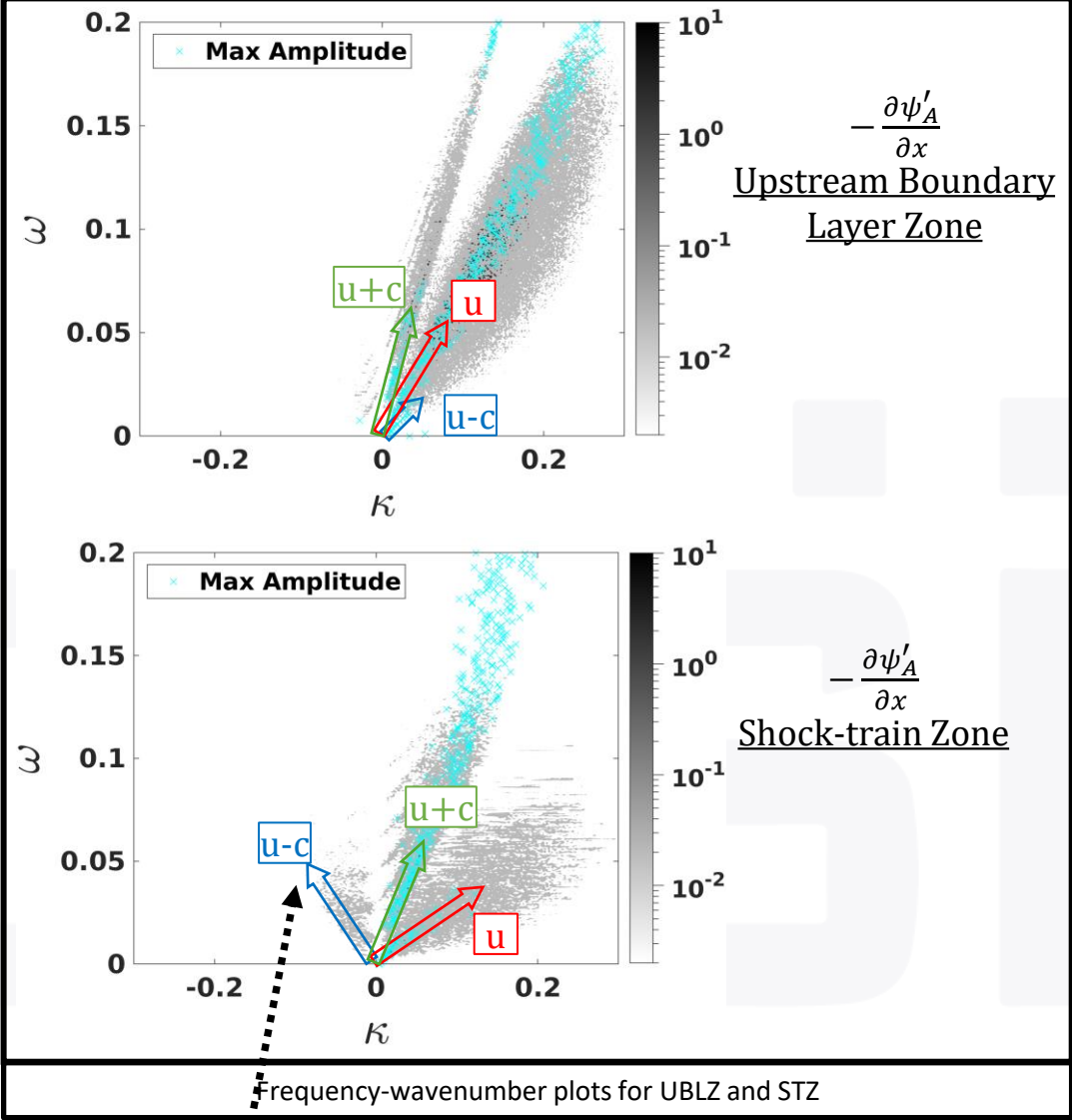
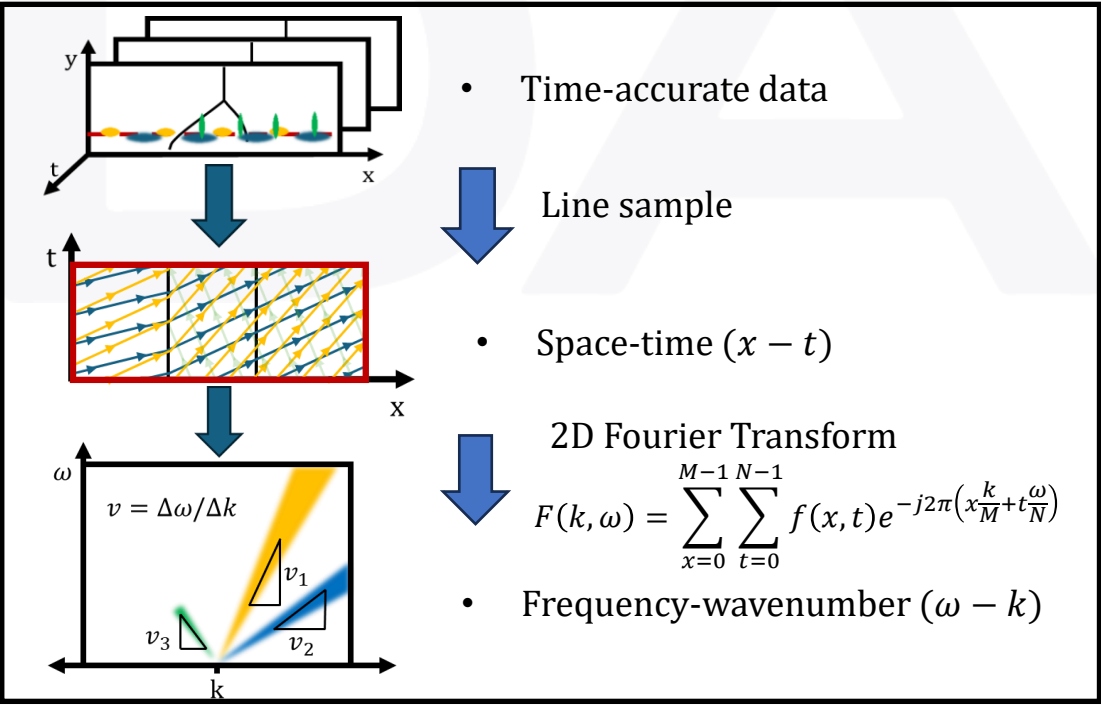
Sample streamwise line near bottom wall and plot over time. Fluctuations only.

2D-Fourier transform into frequency-wavenumber space.

Significance: $v = \Delta\omega/\Delta k$

Positive-slope = downstream | Low-slope = slow

Negative-slope = upstream | High-slope = fast



| | $k - \omega$ | u | $u + c$ | $u - c$ | $c^+ = (u + c) - u$ | $c^- = u - (u - c)$ | c_{avg} |
|------|---------------------------------------|-------|---------|---------|---------------------|---------------------|-----------|
| UBLZ | $-\frac{\partial\psi'_A}{\partial x}$ | 0.710 | 1.641 | 0.384 | 0.931 | 0.326 | 0.629 |
| STZ | $-\frac{\partial\psi'_A}{\partial x}$ | 0.306 | 1.074 | -0.662 | 0.768 | 0.968 | 0.868 |

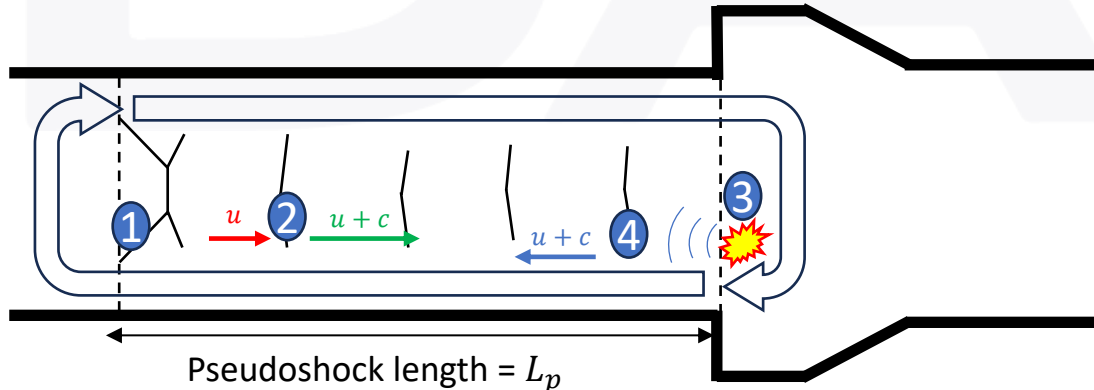
Ensemble averaged upstream-running acoustics

Feedback frequencies

Purpose: Construct scenario for perturbation feedback mechanism within channel and compare frequencies with experiments.

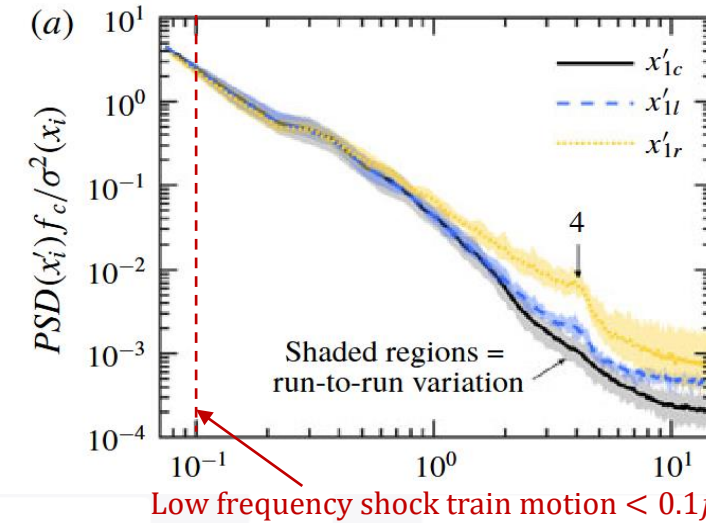
Setup:

1. Assume perturbations begin at 1st shock, purely in streamwise direction
2. Perturbations travel downstream at either u or $u + c$ velocity measured from frequency-wavenumber analysis.
3. End of isolator = beginning of combustor. Assume downstream perturbation creates combustor instability and upstream perturbation.
4. Perturbations travel upstream at $u - c$



Experimental Frequency Spectrum

Hunt and Gamba 2019 [4]



Frequencies are scaled by pseudoshock length.

| | u | $u + c$ | $u - c$ | $c^+ = (u + c) - u$ | $c^- = u - (u - c)$ | c_{avg} |
|--|-------|---------|---------|---------------------|---------------------|-----------|
| $-\frac{\partial \psi'_A}{\partial x}$ | 0.306 | 1.074 | -0.662 | 0.768 | 0.968 | 0.868 |

$$f_{fast} = \frac{1}{\frac{L_p}{u+c} + \frac{L_p}{u-c}}$$

$$f_{fast-scaled} = 0.410f_c$$

$0.410f_c > 0.1f_c$

$$f_{convective} = \frac{1}{\frac{L_p}{u} + \frac{L_p}{u-c}}$$

$$f_{conv-scaled} = 0.209f_c$$

$0.209f_c > 0.1f_c$

Scaled frequencies are **2 to 4 times higher** than the experimental low-frequency unsteadiness

Velocity Filtering

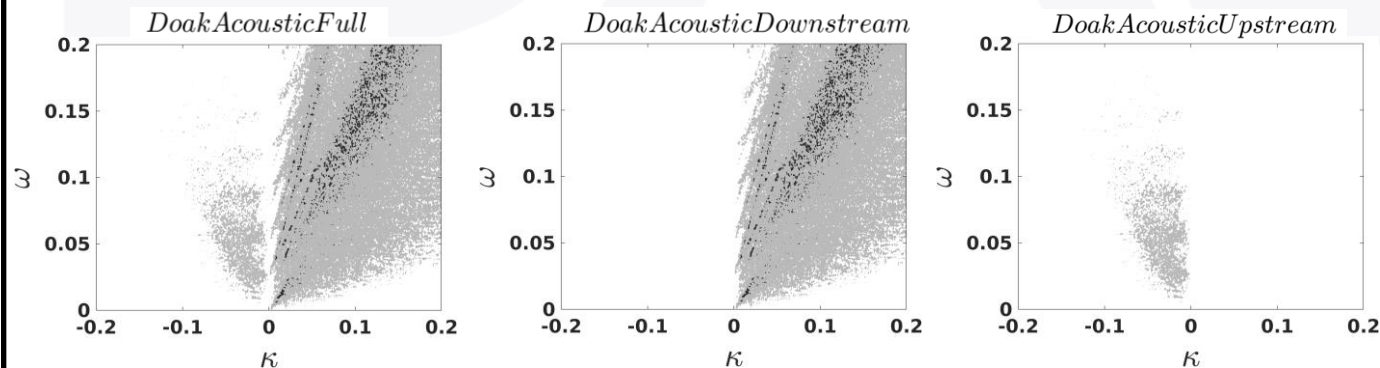
Do perturbations only travel in streamwise direction? Or do they also bounce off walls?

Purpose: Use velocity filtering to isolate upstream and downstream velocity signals in acoustic field.

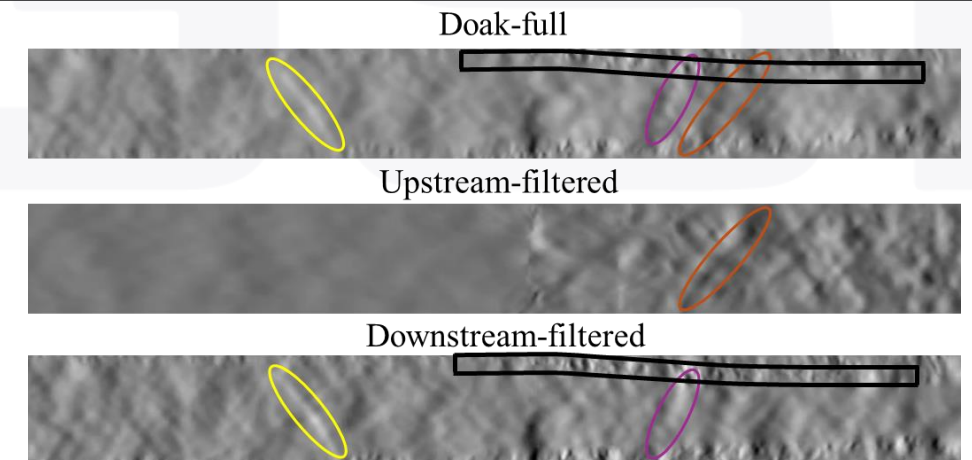
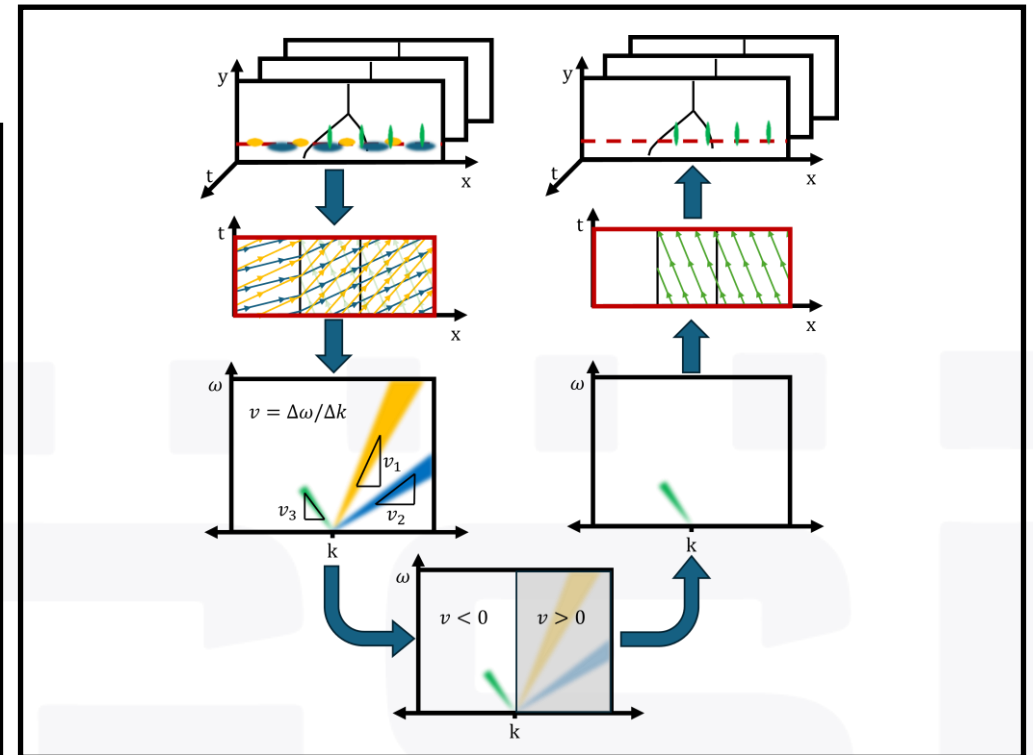
Method: Frequency-wavenumber space can differentiate upstream and downstream waves.

Procedure:

1. Convert line data into frequency-wavenumber space
2. Set positive velocity data to zero
3. Convert back to space-time
4. Repeat for all lines
5. Plot in 2D



Velocity-filtered frequency-wavenumber plots



Waves travel obliquely. Suggests bouncing off walls. May lengthen the pathway of perturbations and reduce frequency

Conclusions

- Time-accurate isolator simulation was constructed.
- Doak's decomposition employed to isolate acoustic perturbations.
- Demonstrated the presence of upstream-running perturbations.
 - First identified using space-time analysis for near-wall streamwise line sample.
- Velocity of perturbations in the shock-train zone were measured.
 - Frequency-wavenumber technique used to obtain ensemble averaged velocities.
- Framework for feedback frequency constructed.
 - Downstream perturbations carried by u and $u + c$ signals.
 - Upstream perturbations carried by $u - c$.
 - Frequencies scaled by pseudo-shock length.
 - Did not agree with experimental low-frequency dynamics.
- Velocity-filtering employed in frequency-wavenumber space.
 - Visually isolating upstream-downstream waves showed angled wave fronts bouncing off walls.
 - Can lengthen propagation pathway, lengthen cycle period, reduce frequency.

References

- [1] Benton, S. I., Stahl, S. L., and Reilly, D., “Characterization of Unsteady Shock Motion in a Transonic Diffuser Flow Path,” AIAA SCITECH 2024 Forum, 2024, p. 1773.
- [2] Dolling D, Or C. Unsteadiness of the shock wave structure in attached and separated compression ramp flows. *Exp Fluids* 1985;3:24–32.
- [3] Ganapathisubramani, B., Clemens, N. T., & Dolling, D. S. “Effects of upstream boundary layer on the unsteadiness of shock-induced separation.” *Journal of fluid Mechanics*, 2007, 585, 369-394.
- [4] Hunt, R. L., and Gamba, M., “On the origin and propagation of perturbations that cause shock train inherent unsteadiness,” *Journal of Fluid Mechanics*, Vol. 861, 2019, pp. 815–859.
- [5] Stahl, S. L., & Benton, S. I., “Computational Investigation of Unsteady Shock Motion in an Isolator-Diffuser Flow Path.” AIAA SCITECH 2025 Forum, 2025, p. 0093.
- [6] Doak, P., “Momentum potential theory of energy flux carried by momentum fluctuations,” *Journal of sound and vibration*, Vol. 131, No. 1, 1989, pp. 67–90.
- [7] Unnikrishnan, S., and Gaitonde, D. V., “Acoustic, hydrodynamic and thermal modes in a supersonic cold jet,” *Journal of Fluid Mechanics*, Vol. 800, 2016, pp. 387–432.